

HELICOIDAL SHELL THEORY

by

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Contents

Contents	5
1 Introduction	7
2 Shell kinematical model	8
2.1 Geometry of the shell	8
2.1.1 Reference systems	8
2.1.2 Oriento-position	8
2.1.3 Base vectors	9
2.1.4 Curvature.....	9
2.2 Constant curvature motion of the material surface	10
2.2.1 Geometry of the material surface	10
2.2.2 Constant rototranslation speed	10
2.2.3 Tangent curvature dual vectors	10
2.2.4 Oriento-position speed dual vector.....	11
2.2.5 Constant curvature motion	11
2.3 The helicoidal shell model	12
2.3.1 The oriento-position field.....	12
2.3.2 Model curvatures.....	13
2.3.3 Model base vectors.....	14
2.3.4 Model curvature tensor.....	15
2.3.5 Model surface and volume elements	15
2.3.6 Relaxation of the kinematical hypotheses	16
2.3.7 Initialization of the dual director variable	17
2.3.8 Geometric invariance	17
2.4 Deformed configuration	17
2.4.1 Deformed model oriento-position	18
2.4.2 Deformed model curvature.....	18
2.4.3 Deformed model base vectors and deformation gradient	19
2.5 Rototranslation of the shell model	19
2.5.1 The rototranslation field	19
2.5.2 Shell model strain.....	20
2.5.3 The kinematical strain	20
2.6 Perturbed configuration	21
2.6.1 Co-rototranslational variations of the model kinematical strain	21
3 Shell model mechanics	22
3.1 Shell model equations	22
3.2 Variational formulation	23
3.2.1 Shell model variational framework	23
3.2.2 Boundary terms and boundary constraints	24
3.2.3 Principle of Virtual Work.....	25
3.3 Consistent linearization	25
3.3.1 Internal stress contribution	25

4 Shell model local constitutive law.....	27
4.1 Nonlinear elasticity	27
4.1.1 Transverse strain local variable	28
4.1.2 Transverse strain local condensation	31
4.1.3 Transverse kinematical strain restraint	33
4.2 Linear elasticity.....	35
4.3 Linearized virtual functional.....	37
5 Reduction to the shell surface mechanics	38
5.1 Shell Biot-axial model.....	38
5.1.1 Substitute Biot-axial field.....	38
5.1.2 Biot-axial field variations	39
5.2 Surface density of the internal stress virtual functional	39
5.2.1 Shell model internal stress virtual functional	39
5.2.2 Material surface internal stress virtual functional.....	41
5.2.3 Recovering the transverse strain local variable	42
5.3 Shell constitutive mechanism.....	42
5.3.1 Proposition	42
5.3.2 Linearized Principle of Virtual Work of the shell	43
5.3.3 Shell director and Biot-axial local variables.....	44
5.3.4 Shell constitutive equations and tangent map.....	45
5.3.5 Alternative choices of the Biot-axial local variables	48
Appendices	51
A1 Kinematical strain variations	52
A1.1 Curvature co-rototranslational variation formulae.....	52
A1.1.1 Setting the co-rototranslational variation formulae.....	52
A1.1.2 Reducing the co-rototranslational variation formulae.....	56
A1.2 Kinematical strain co-rototranslational variation formulae	61
A2 Hyperelastic constitutive laws	63
A2.1 Three-dimensional models for non-polar isotropic materials	63
A2.1.1 The Saint Venant-Kirchhoff constitutive model.....	63
A2.1.2 The Neo-Hookean constitutive model	64
A2.1.3 Saint Venant-Kirchhoff constitutive model by Biot-type parameterization.....	65
A2.1.4 A linear constitutive law by Biot-type parameterization	66
A2.1.5 Neo-Hookean constitutive model by Biot-type parameterization	67
A2.1.6 Pseudo-Hookean constitutive model by Biot-type parameterization	68
A2.1.7 Synopsis	70
A2.2 Shell model local constitutive laws for non-polar isotropic materials	71
A2.2.1 Nonlinear constitutive laws	71
A2.2.2 Linear constitutive law	72
A2.2.3 Linear constitutive law using technical constants.....	76
A3 Internal stress virtual functional.....	80
A3.1 Surface density of the internal stress virtual functional	80
A3.1.1 Shell model internal stress virtual functional.....	80
A3.1.2 Material surface internal stress virtual functional.....	87
A3.1.3 Recovering the transverse strain local variable.....	88
References	91

1 Introduction

The present report attains the second step in the undertaking of introducing the *helicoidal modeling* (Merlini and Morandini, 2004a, 2004c) in computational shell mechanics. The main variational formulations of the *shell material surface* in the context of the helicoidal modeling have been developed in a first report by Merlini (2008b). That report focused on the intrinsic mechanics of the shell described by a two-coordinate manifold, whereas in the present second report, we proceed to recover the *solid shell* constitutive mechanism from three-dimensional continuum mechanics within a helicoidal modeling context.

The theory presented herein is composed of the solid shell kinematical model and the relevant variational mechanics. The shell is regarded as a three-dimensional continuum made of infinitesimal yet three-dimensional material particles, susceptible of a polar description. A kinematical model of this solid shell is a substitute three-dimensional body made of exactly the same material particles, but whose oriento-positions are described as known functions of a discrete set of parameters across the thickness (a kind of approximation close to the finite element method). The kinematical model we propose stands on very simple hypotheses in the context of the helicoidal modeling of the continuum: *the curvature dual vectors*, that characterize the oriento-position gradient, *are assumed constant across the shell thickness*. As a result, the model configuration depends on the configuration of the underlying material surface (Merlini, 2008b), and on one *dual director* that characterizes the generalized curvature across the thickness.

The variational mechanics of the solid shell model is nothing but an application of the variational formulation for three-dimensional continuum mechanics (Merlini and Morandini, 2004c, 2005). For simplicity, only a hyperelastic non-polar medium is considered, and just the ‘one-field’ formulation – the internally constrained Principle of Virtual Work – is written and consistently linearized. To represent the axial of the Biot stress within the shell model, a suitable and simple substitute field is hypothesized: *the local Biot-axial is assumed linear across the thickness* and depending on two parameters that are conveniently connected to the *dual Biot-axial* field of the underlying material surface (Merlini, 2008b). Once the kinematical and the Biot-axial models are substituted within the variational formulation and the appropriate integrals have been computed across the thickness, the Principle of Virtual Work must correspond to that obtained for the material surface. A further unknown is present now, however, besides the oriento-position and dual Biot-axial of the material surface: the dual director of the shell model.

The shell theory so far outlined yields the constitutive law to exploit in the variational mechanics of the material surface. The constitutive mechanism descends from an obvious statement of energetic equivalence between the virtual internal work formulated for the solid shell model and that of the material surface. This statement stands on a nonlinear virtual functional and is extended to its linearized form, which is continuously rebuilt during the incremental solution process. Since the dual director unknown does not pertain to the material surface mechanics, it would be convenient to solve it within the constitutive characterization itself, by means of a local condensation technique dynamically operated on the linearized tangent map. This condensation process needs a constraint equation and an unknown reaction to be accomplished, and the angular part of the surface dual Biot-axial is used for this purpose, leaving a surface mechanics theory endowed with only the oriento-position and the linear part of the surface dual Biot-axial as unknowns. The result of the constitutive mechanism is the expression of the stress resultants (forces and couples per unit length) and the relevant derivatives with respect to the kinematical strains of the material surface, namely the shell *constitutive equations* and their *tangent* map. Such quantities are integrals over the shell thickness, numerically obtained during the iterative solution process.

The Report is organized as follows. In Section 2 we introduce and discuss the constant-curvatures shell kinematical model. In Section 3 we reformulate the three-dimensional variational mechanics for the solid shell model and state and linearize the Principle of Virtual Work. In Section 4 we accommodate the local constitutive law for the particular shell model kinematics. In Section 5 we focus on the internal work linearized functional: we propose a substitute Biot-axial field and perform the integrals over the shell thickness; then, the shell constitutive equations and their tangent map are formulated. In the appendices we give details of the derivation of the kinematical strain variations (Appendix A1), of some basic local hyperelastic models (Appendix A2), and of the linearization of the internal stress virtual functionals (Appendix A3).

Notation

Throughout this Report, Latin indexes denote generic tensor components in the three-dimensional space, so they are intended to vary in the range 1 to 3. Instead, Greek indexes denote components relating to the shell two-dimensional surface and are intended to vary in the range 1 to 2.

The Einsteinian rule of implicit summation over repeated indexes is usually understood.

2 Shell kinematical model

As *solid shell*, we understand a three-dimensional body, lying on a smooth curved surface and thin in the direction locally normal to the surface. The thickness of this three-dimensional shell body is not necessarily constant over the surface; however, it is allowed to vary in a quite smooth manner so that the solid shell can always be referred to as a thin solid.

In agreement with the polar description of material solids, the solid shell is defined as a continuum set of infinitesimal yet three-dimensional material particles occupying the volume of a three-dimensional shell body. The status of each particle is identified by its position and by the orientation of an embedded triad of independent vectors. In a deformable shell, the material particles may change their relative distances and rotations, yet they keep on belonging to a deformed three-dimensional shell body, again lying on a smooth curved surface and thin in the direction locally normal to the surface.

In this chapter, we propose a kinematical model of the solid shell based on the *helicoidal modeling* of the continuum. Although this model is an approximation of the actual shell kinematics, it provides a natural description — suitable for curved thin bodies — that can be advantageously exploited in computational shell mechanics. The description of the solid shell is formulated here as an extension of the *material surface* description, as established in Merlini (2008b). The reader should refer to Chapter 2 of that Report as the fundament of the present shell kinematical model.

2.1 Geometry of the shell

The geometry of the solid shell is formulated in exactly the same way as for any three-dimensional solid (Merlini and Morandini, 2002, 2004c). In this section, we recall the essential steps of the formulation in the context of the helicoidal modeling.

2.1.1 Reference systems

First, we assume an absolute reference frame. We take as *absolute frame* an orthonormal triad of dimensionless unit vectors $\mathbf{i}_j \equiv \mathbf{i}^j$, fixed and located somewhere in the three-dimensional space, at a point called the absolute origin. This absolute frame, though arbitrary, is completely known, in both position and orientation, to anybody. The absolute frame is represented by the identity real tensor $\mathbf{I} = \mathbf{i}_j \otimes \mathbf{i}^j$.

Associated with this frame, we conveniently define an absolute system of coordinates. The *absolute reference system* is based on Cartesian coordinates measured along three orthogonal axes x^j issued from the absolute frame at the origin and oriented as the unit vectors \mathbf{i}_j , respectively. This reference system is implicitly assumed throughout the present theory; in fact, for numerical purposes, every variable in a formulation of continuum mechanics shall be ultimately measured in a unique, absolute system.

For the theoretical developments, we also resort to an intrinsic reference system made of convective coordinates. This *material reference system* is based on (dimensionless) curvilinear coordinates measured along three families of smooth lines ξ^j traced within the three-dimensional shell body. As usual, the coordinates ξ^α ($\alpha = 1, 2$) are assumed to lie on surfaces ‘parallel’ to the shell faces, in the sense that they never cross the lower and upper faces of the solid shell. Instead, the coordinate lines ξ^3 do cross them. The coordinate surfaces (ξ^1, ξ^2) will be often referred to as *layers*. The lower and upper shell faces are coordinate surfaces themselves, located at some abscissas ξ^{3-} and ξ^{3+} , respectively.

2.1.2 Oriento-position

A material particle of the solid shell is identified by its position vector \mathbf{x} and by the orientation tensor $\boldsymbol{\alpha} = \boldsymbol{\alpha}_j \otimes \mathbf{i}^j$ of a triad of vectors $\boldsymbol{\alpha}_j$ embedded in it. This triad must be unique for all the particles except for a relative rotation, so the particle orientations can also be understood as rotations from a unique triad. Although it is not strictly necessary, we assume that this unique triad is the absolute orthonormal frame itself: by doing so, the orientation tensors $\boldsymbol{\alpha}$ are actually rotation tensors.

A dual frame, as defined in Merlini and Morandini (2004c), is a complex geometrical entity suited to identify at the same time the position and the orientation of a material particle. This leads to a pole-based description and entails the choice of a pole. If we make the pole to coincide with the origin itself, a material particle is identified by the origin-based *oriento-position* dual tensor